- 1) Gas adsorption studies, and
- 2) Motions of ions and molecules in crystals and on nanostructured surfaces

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- Gas Physisorption nanostructured surfaces, gases at high pressure technical comments about measurements
- 2) Jump Mechanisms use pressure to determine V_A transient structure for the jump polarons, ions, and molecules

Sorption Methods for Hydrogen Storage

ABsorption (chemisorption)



Something new in bcc VTiNiCr

ADsorption (physisorption)



Graphene from graphite oxide Nanostructured surfaces

Adsorption Isotherms f(T,P)

For a fixed temperature, the probability that a surface site is occupied is

$$f_{\rm ad} = \frac{\mathrm{e}^{(\mu-\varepsilon)/k_{\rm B}T}}{1 + \mathrm{e}^{(\mu-\varepsilon)/k_{\rm B}T}} , \qquad (16.26)$$

two possibilities -

(

empty site (with Gibbs factor $e^0 = 1$) site with an adsorbed molecule (with Gibbs factor $e^{(\mu - \epsilon)/k_BT}$)

Equilibrium: chemical potential of adsorbed molecule equals that of a molecule in the gas. Get chemical potential of ideal gas:

$$e^{\mu/k_{B}T} = \frac{N V_{Q}}{V}$$

$$pV = Nk_{B}T$$

$$f_{ad} = \frac{p}{\frac{k_{B}T}{V_{Q}}}e^{\epsilon/k_{B}T} + p$$

$$p_{0}(T) \equiv \frac{k_{B}T}{V_{Q}}e^{\epsilon/k_{B}T}$$

$$f_{ad} = \frac{p}{p_{0}(T) + p}$$

$$f_{\rm ad} = \frac{p}{p_0(T) + p}$$



Tuning Physisorption Thermodynamics with Nanostructured Surfaces



A porous material surface enhanced non-ideal gas interactions between methane molecules at high pressures.

Adsorbent Properties





Methane Intermolecular Interactions on Nanostructured Surface



4.2Å

Methane Shows Anomalous Isosteric Heat for ZTC



Together Methane and Ethane account for ~98% of natural gas.

Ethane Data Fitted with Superposition of Langmuir Isotherms



Fit data with:

$$n_e(P,T) = (n_{max} - V_{max}\rho(P,T)) \left[\sum_i \alpha_i \frac{K_i P}{1 + K_i P}\right]$$

Isosteric Heat of Ethane Adsorption



Calculated with the Clapeyron Equation:

$$q_{st} = -\Delta H_{ads}(n_a) = -T\left(\frac{\partial P}{\partial T}\right)_{n_a} (\Delta v_{ads})$$

Entropy of the Adsorbed Phase (Ethane)



Ongoing Work: Krypton Adsorption



Volumetric Adsorption Measurements



Sieverts method:

$$\rho(P_i,T) \; V_i = \rho\bigl(P_f,T\bigr) \; V_f + n_{ads}$$

(whole apparatus at uniform temperature)



For Samples of ~1000 m² signal to noise \approx 100:1 For Sample of ~100 m² signal to noise \approx 10:1 100 m² is minimum for adsorption measurements

Activation Volumes of Kinetic Mechanisms





Classical: Ion or atom Li+ senses only initial state and transition state Quantum: Polaron Electron (hole) samples both initial and final states.

Measure $I(E,P,T) \rightarrow \Gamma(P,T) \rightarrow V_A$ Get Activation Volume Activation volume gives a picture of atoms in transition state

Activation Volume in Li_xFePO₄

Originates from the local atomic distortions that bring the electron levels of the initial and final states into coincidence, facilitating electron transfer

$$\Gamma = \Gamma_0 \exp(-2\alpha R) \exp\left(-\frac{E_a + PV_a}{kT}\right)$$

Is V>0 or V<0?

 $V_{\rm a}$ is determined by measuring the effect of applied pressure on the polaron hopping frequency

Electrical conductivity of oxides showed: $V_a = -0.2 \text{ Å}^3$ We found: $V_a = +6 \text{ Å}^3$

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Dynamical Effects in Mössbauer Spectra

•

$$\tau_{M} = \frac{\hbar}{E_{HF}}$$
• Observe ⁵⁷Fe spectral
relaxations from Fe²⁺–Fe³⁺
• Account for the time
correlations of a fluctuating
Hamiltonian with Blume-
Tjon model.
$$(velocity)$$

$$\tau - \tau_{M}$$

Nuclear Forward Scattering at APS 16ID-D (HPCAT)

DAC in Copper Block Furnace







Heating HPCAT Block For Symmetric DAC

Uniform heating up to 500C

Good temperature stability

Integrated membrane pressure control

Remote Temperature & Pressure Control

Front plate has access holes for the symmetric cell bolts

back side plate constrains a membrane for remote pressure control





Li_xFePO₄ Nuclear Forward Scattering

Spectra at 0 GPa show large distortions, consistent with the polaron dynamics reported by conventional Mössbauer

At elevated pressure these large spectral distortions do not occur until higher temperature

Spectral Distortions appear to be more pronounced at low Lithium concentration



Estimates of Concentration Dependence of Activation Volume



 $\begin{array}{l} {\rm Li}_{.4}{\rm FePO}_4 \ {\rm V}_{\rm A} \sim 9 {\rm \AA}^3 \\ {\rm Li}_{.66}{\rm FePO}_4 \ {\rm V}_{\rm A} \sim 6 {\rm \AA}^3 \\ {\rm Li}_{.8}{\rm FePO}_4 \ {\rm V}_{\rm A} \sim 3 {\rm \AA}^3 \end{array}$

Activation volume has a roughly linear increase as lithium concentration is reduced





Compressibility and Activation Volume

Destabilization of an electron polaron centered at a Fe²⁺ ion because the compressibility of ferrous-oxygen bonds is greater than for ferric-oxygen bonds

The Fe³⁺-oxygen bonds are shorter and stiffer than the Fe²⁺-oxygen bonds \rightarrow we expect a larger the effect of pressure on the ferrous Fe-O bonds

Overall reduction of activation volume should be greater at low Li concentrations

Opposite trend seen in experimental results \rightarrow not dominant effect

Perhaps the dominant effect is how pressure alters the activation volume for Li⁺ ions in materials with higher Li concentration?

Quasielastic Scattering (Neutrons and Hydrogen)

Incoherent scattering, analysis of diffusion is good (diffraction is not).

¹H has neutron scattering cross section σ ~80 barns

(10 barns is big for other isotopes)

For inelastic scattering, σ/m of ¹H is unbeatable

Neutron scattering from pressure cells can be less than from ¹H



FIG. 1. Gasket-anvil assembly for high pressure QENS measurements. (1) and (2) anvil made of sintered cBN die supported by steel binding ring; (3) split CuBe gasket; (4) high pressure sample chamber; (5) cartridge heater; (6) and (7) anvil backing seat made of zirconia die and steel supporting ring; (8) cooling rings. Conical anvil faces (thick lines) are covered with gadolinium film.

Bove, et al, "Translational and Rotational Diffusion in Water in the Gigapascal Range," PRL 111, 185901 (2013).

Quasielastic Scattering (General)

Van Hove function *G*(*r*,*t*)

$$\begin{split} G(\vec{r},t) &\equiv \frac{1}{\tau} \int_{t'=0}^{+\tau} \int_{\vec{r}=-\infty}^{+\infty} f^*(\vec{r}'',t') f(\vec{r}+\vec{r}',t+t') \,\mathrm{d}^3 \vec{r}' \,\mathrm{d}t' \\ S(\vec{Q},\omega) &= \frac{1}{2\pi\hbar} \int_{-\infty}^{+\infty} G(\vec{r},t) \,\mathrm{e}^{-\mathrm{i}(\vec{Q}\cdot\vec{r}-\omega t)} \,\mathrm{d}^3 \vec{r} \,\mathrm{d}t \\ S(\vec{Q},\omega) &= \frac{1}{2\pi\hbar} F_{\vec{r}} F_t G(\vec{r},t) \,, \end{split}$$

Calculate the Van Hove function for diffusion (different atoms are uncorrelated)

$$c(x,t) = \frac{c'}{\sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}} \qquad G_{\text{diff}}(x,t) = \int_{-\infty}^{+\infty} c(x',0) c(x'+x,t) \, \mathrm{d}x'$$

$$S(Q,\omega) = \frac{c'}{2\pi\hbar} \frac{DQ^2}{(DQ^2)^2 + \omega^2} \qquad \Delta E = \hbar D Q^2$$

Quasielastic Scattering and Pressure



FIG. 3 (color online). Example of fit result (red line) compared to the measured intensity (blue dots) from water at 400 K at P = 0.6 and 1.5 GPa for *Q*-selected values.

Quasielastic Scattering and Pressure

Klotz, et al., "Quasi-elastic neutron scattering in the multi-GPa range and its application to liquid water", APL 103, 193504 (2013).



FIG. 3. Full width at half maximum (FWHM) of the fitted QENS signal plotted as a function of Q^2 , at 0.6 GPa and 2.9 GPa. The marked dots (\odot) at 0.86 Å⁻² are those shown in Figure 2. Full symbols correspond to analysis

Activation volume:

My estimate is $+2 \text{ Å}^3$ for translational diffusion of water molecule in liquid. (*T*-Dependence?)

Activation Volume for ¹H Diffusion

- Activation volume is of interest for both physisorption and chemisorption.
- Plans at the SNS to do inelastic scattering under pressure.
- If the activation volumes are several Å³, pressure of only 1 GPa are required for measuring V_A .
- Coordination with the SNS needs thought.